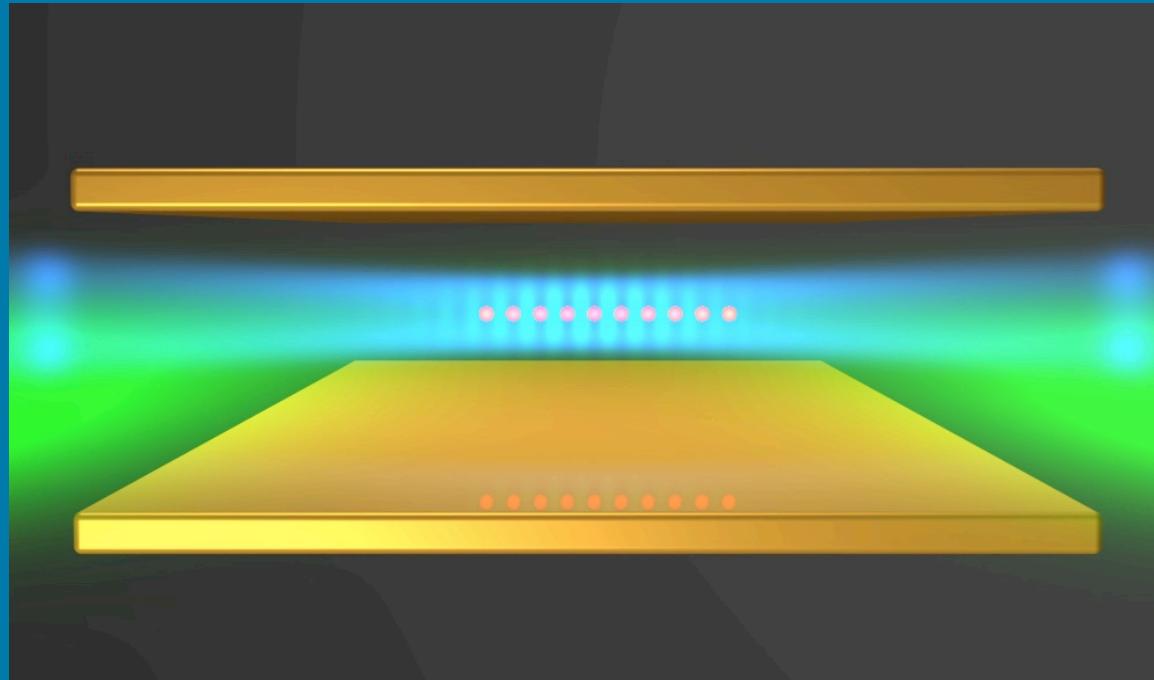


# Towards Quantum Simulation with Circular Rydberg Atoms

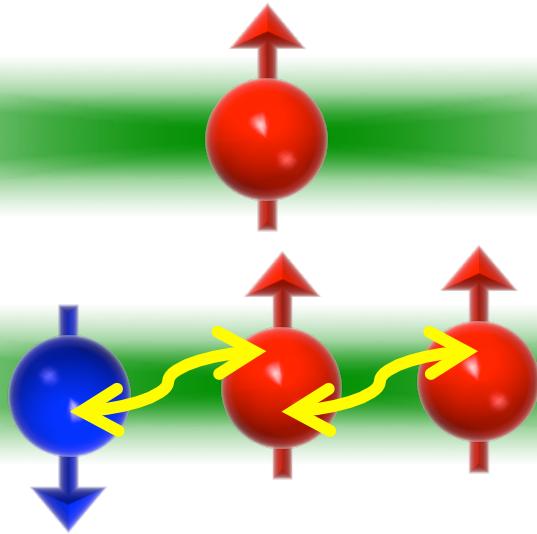


Clément Sayrin  
Laboratoire Kastler Brossel

ICQSIM, 15/11/2017

# Quantum simulations of spin-systems

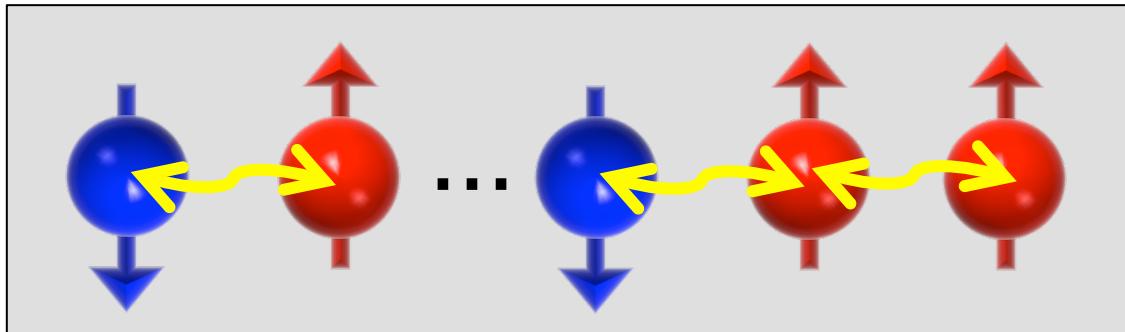
Simulations with computer?



Simple

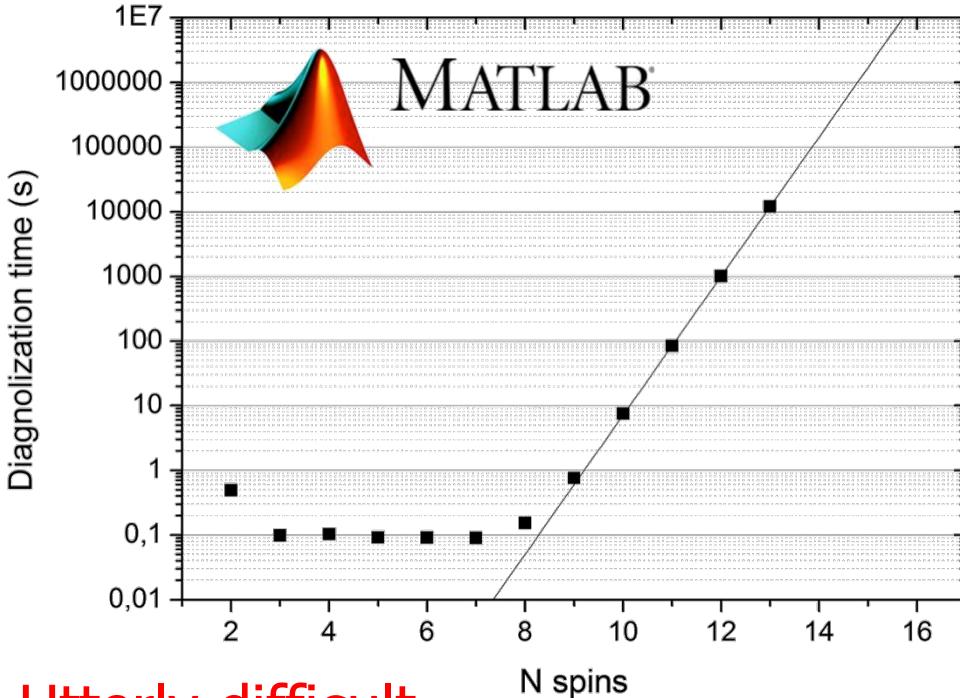
Barely more complicated

Utterly difficult

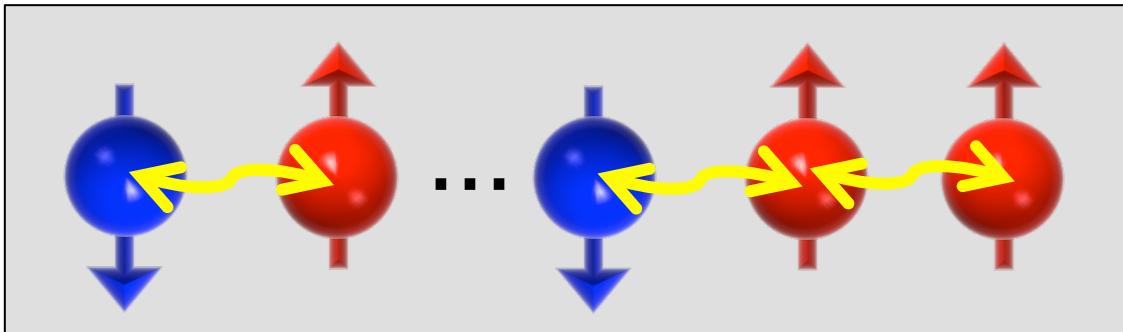


# Quantum simulations of spin-systems

Simulations with computer?



Utterly difficult



$N = 10 \rightarrow \sim 7 \text{ s}$

$N = 13 \rightarrow \sim 3 \text{ h}$

$N = 16 \rightarrow \sim 8 \text{ months}$

$N = 40 \rightarrow \sim 4 \cdot 10^{25} \text{ years}$



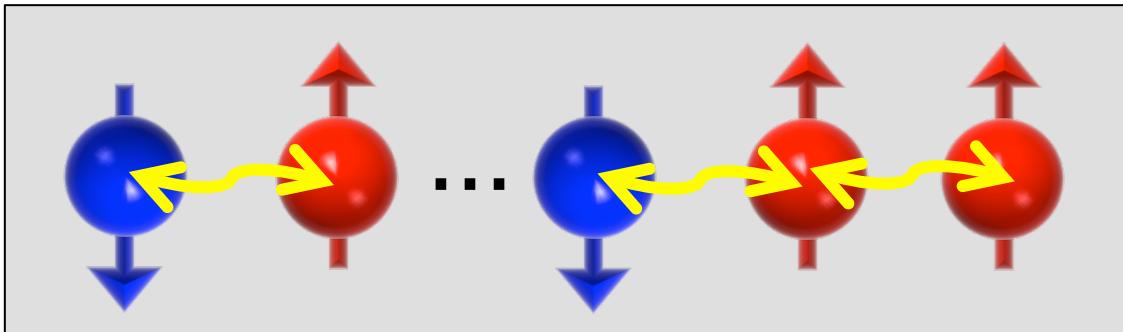
# Quantum simulations

Simulations with computer?

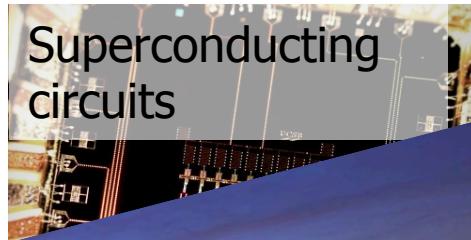
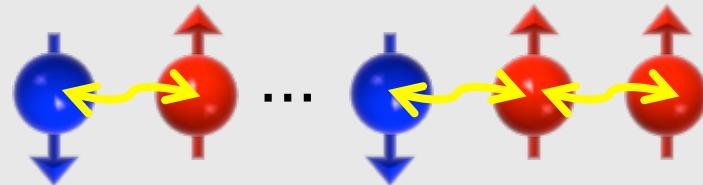
Linear chain of  $N$  spins-1/2:

- Exact diagonalization:  $N \sim 36$
- Ground-state: well-known via powerful numerical techniques (DMRG)
- Dynamics... tricky! Few tens of interaction cycles only.

Utterly difficult



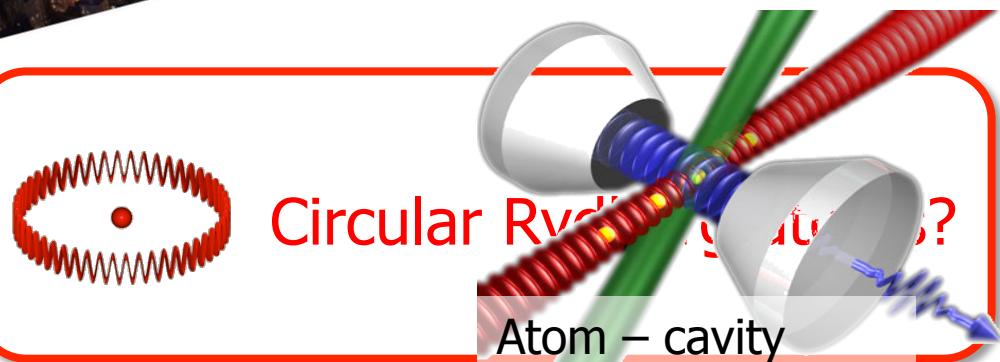
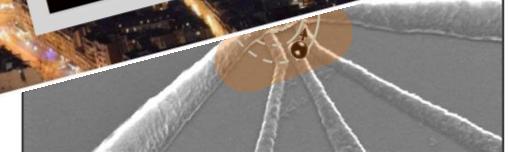
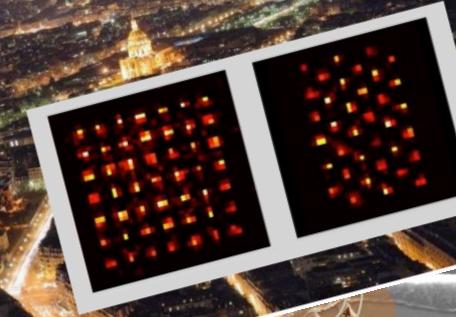
# Quantum simulations



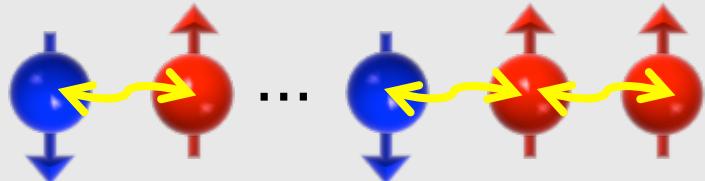
How to simulate?

Use a model

ICQSIM – 2017  
International Conference on  
Quantum Simulation  
13-17 Nov. 2017 – Paris (France)



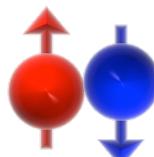
# Objective



Simulation of a chain of interacting spins-1/2

## Requirements?

- Spin 1/2



- Defect free chain of spins
- Long lifetime and strong interaction



Observe many interaction cycles

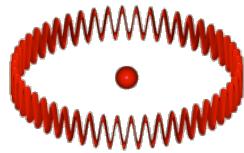
- Fully tunable Hamiltonian  $H = H_0 + H_{\text{ext}} + H_{\text{int}}$



# Circular Rydberg atoms $|nC\rangle$

## Rydberg atoms

- Very high principal quantum number,  $n$
- Circular levels: maximum angular momentum  $l = |m| = n-1$



Large orbit  $r_n \sim n^2 a_0$

→ Huge electric dipole matrix elements

- Well coupled to the microwave electromagnetic field
- Strong dipole – dipole interactions!

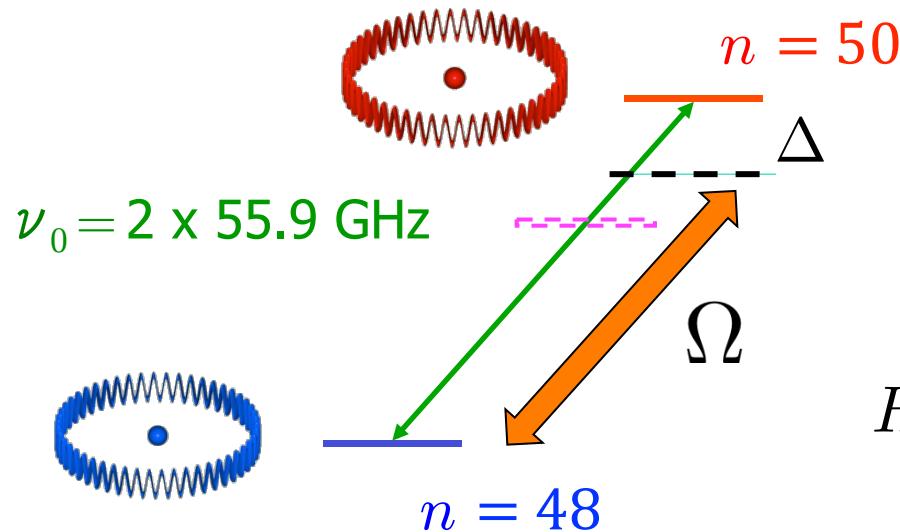
+ Long lifetimes: Several 100 $\mu$ s for low  $l$  levels

Several 10 ms for circular levels

# Circular Rydberg atoms

## Simulating a spin 1/2

- Two circular Rydberg levels  $|nC\rangle$   $|(n + 2)C\rangle$  + near-resonant drive



$$H_0 + H_{\text{ext}} = \frac{h\nu_0}{2}\sigma^z + \frac{h\Delta}{2}\sigma^z + \frac{h\Omega}{2}\sigma^x$$

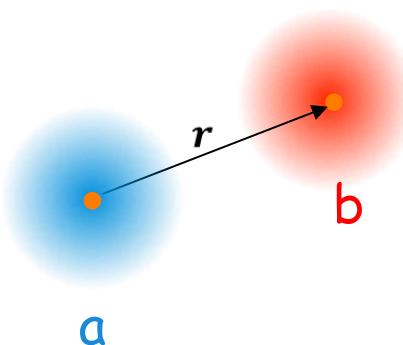
Longitudinal magnetic field

Transverse magnetic field

Power and frequency of the microwave source



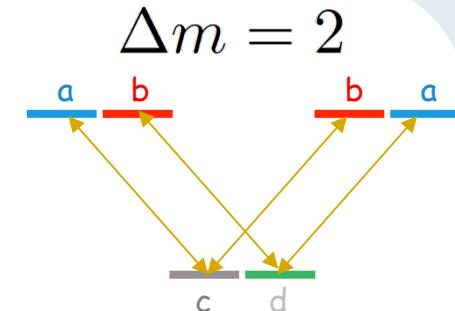
# Dipole-dipole interaction



## Off-diagonal terms



Exchange  
interaction



$$\Delta m = 2$$

$$H_{\text{int}} = \sim \hat{d}_1 \hat{d}_2 / r^3$$

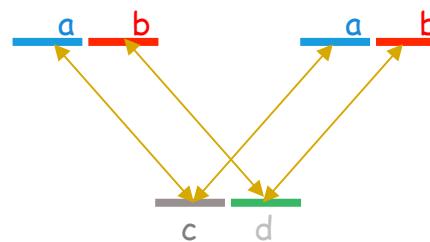
$|ab\rangle$   $|ba\rangle$

$C_{a,b}$   $A_{a,b}$

$A_{b,a}$   $C_{a,b}$

A diagram showing the interaction Hamiltonian  $H_{\text{int}}$  proportional to  $\hat{d}_1 \hat{d}_2 / r^3$ . It shows two states:  $|ab\rangle$  and  $|ba\rangle$ . Between these states are four interaction terms:  $C_{a,b}$  (blue),  $A_{a,b}$  (orange),  $A_{b,a}$  (orange), and  $C_{a,b}$  (blue).

## Diagonal terms



Direct  
interaction

Second order, van der Waals  
interactions

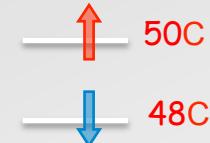
$$J_z, J \propto 1/r^6$$

Nearest-neighbour interactions

## Mapping to a spin interaction

Diagonal terms

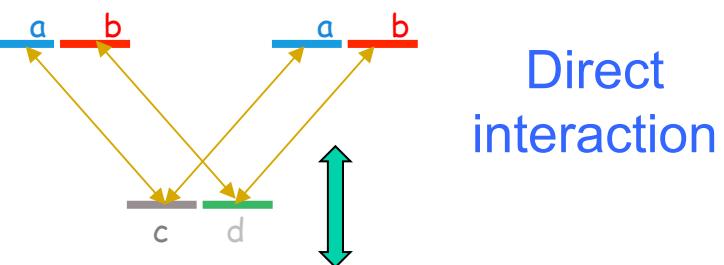
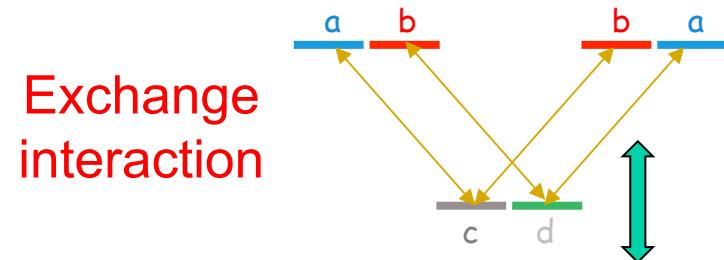
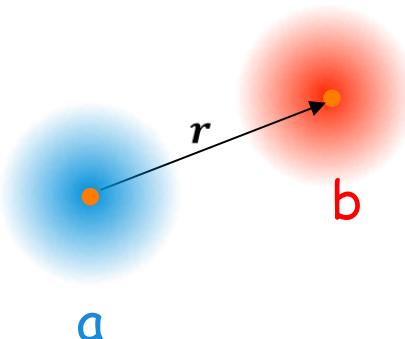
$$J_z \sigma_1^z \sigma_2^z$$



Off-diagonal terms

$$J (\sigma_1^x \sigma_2^x + \sigma_1^y \sigma_2^y)$$

# Tunable spin – spin interaction



- Tunable dipole – dipole interaction:  
 $A_6$  and  $C_6$  coefficients depend on  $F$ ,  $B$
- Tunable spin – spin interaction:

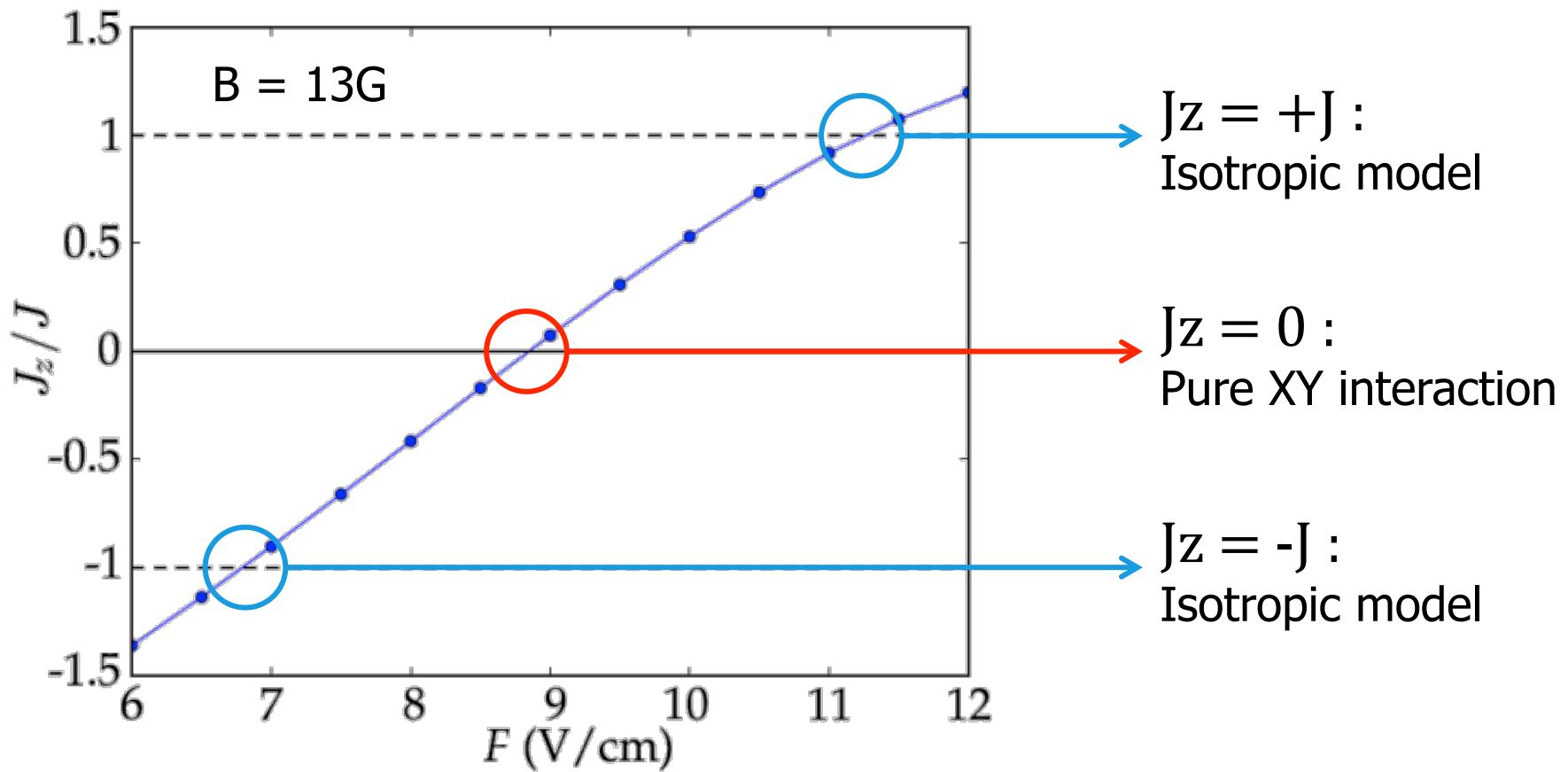
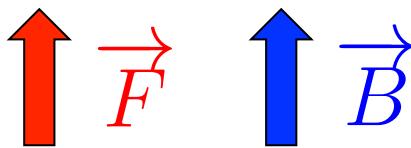
$$J = \frac{A_{6,48-50}}{2r^6}$$

$$J_z = \frac{C_{6,48-48} - 2C_{6,48-50} + C_{6,50-50}}{4r^6}$$

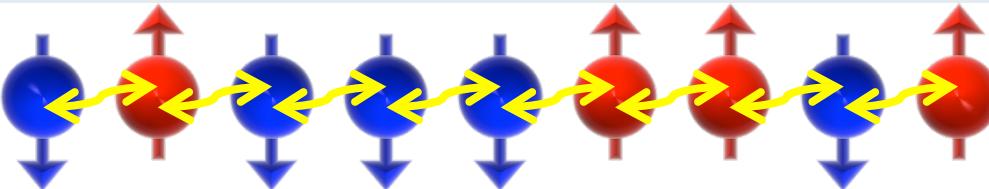
$J$  is nearly constant...

... but  $J_z$  varies significantly!

# Tunable spin – spin interaction



# Tunable XXZ Hamiltonian



$$H/h = J_z \sum_j \sigma_j^z \sigma_{j+1}^z + J \sum_j (\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y)$$

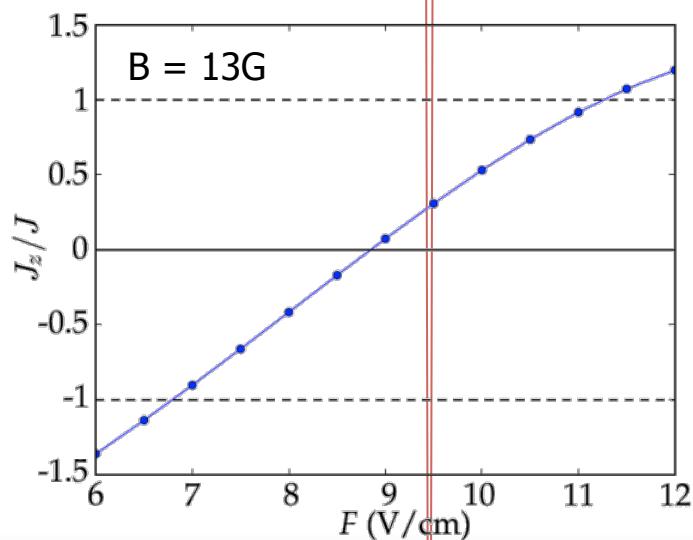
Longitudinal B field      Transverse B field

$$\frac{\Delta}{2} \sum_j \sigma_j^z$$

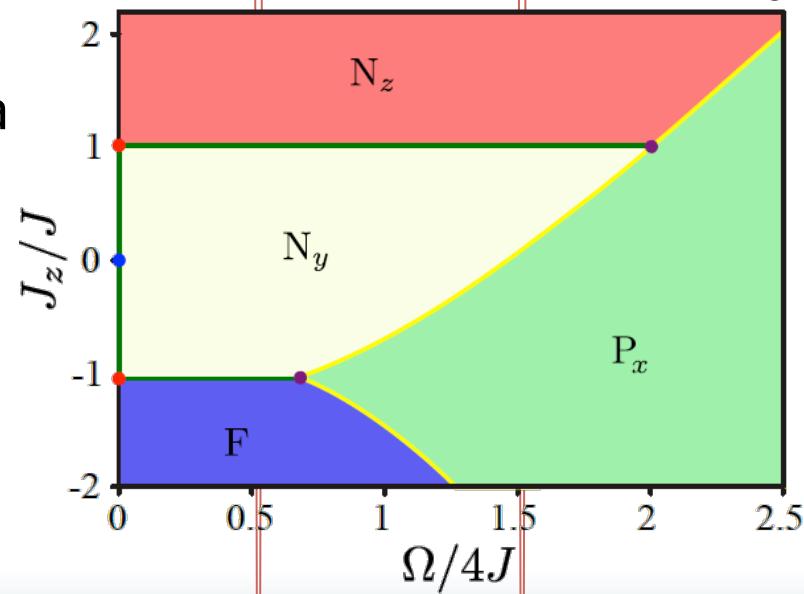
$$\frac{\Omega}{2} \sum_j \sigma_j^x$$

$$\Delta = 0$$

## Tunable interactions



Exploration of a rich phase diagram

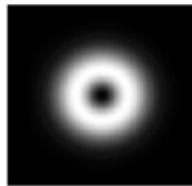


# Trapping Rydberg atoms

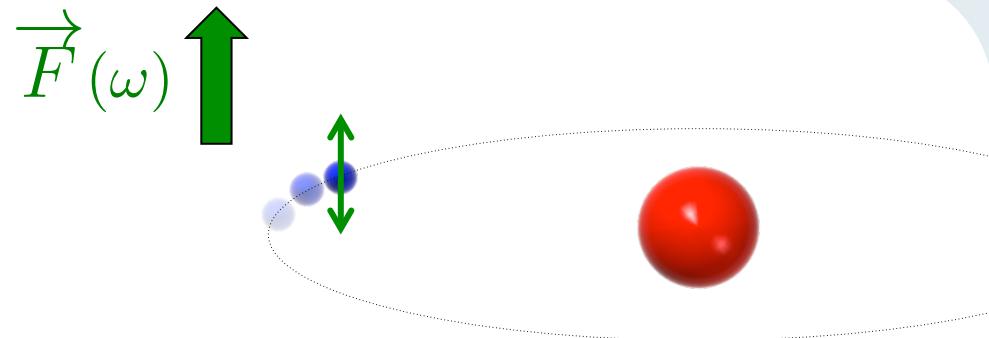
## Ponderomotive potential

- Valence electron is almost free
- Positive ponderomotive energy

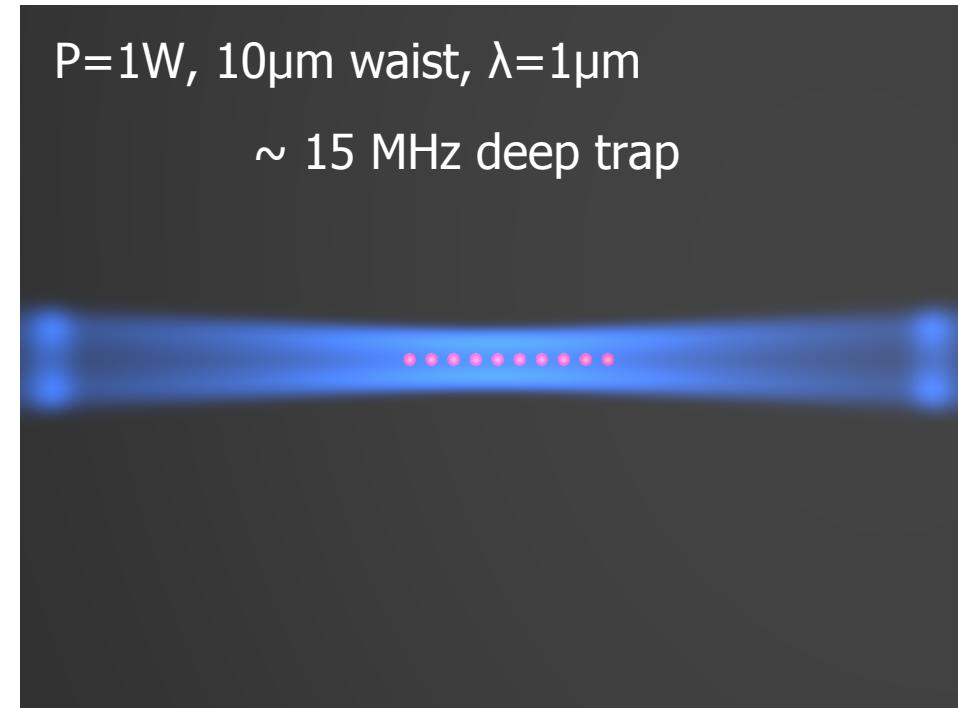
$$\mathcal{E} = \frac{e^2}{2m_e \varepsilon_0 c \omega_L^2} I \quad \text{Low field seeker}$$



- Laguerre-Gauss beam
- Transverse confinement



P=1W, 10µm waist, λ=1µm  
~ 15 MHz deep trap

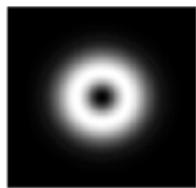


# Trapping Rydberg atoms

## Ponderomotive potential

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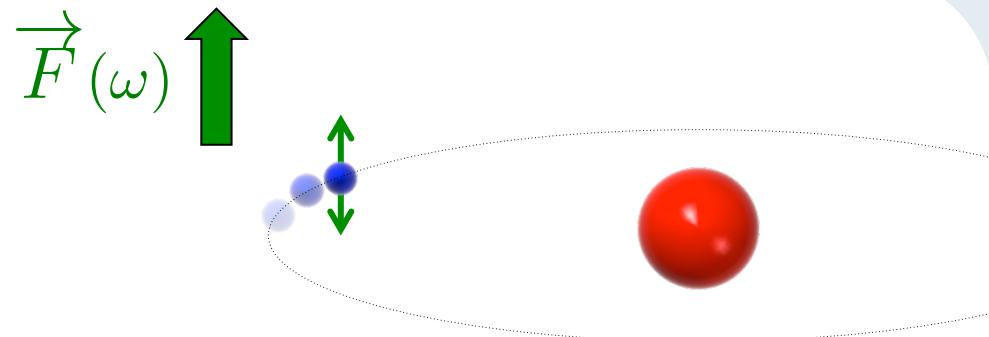
$$\mathcal{E} = \frac{e^2}{2m_e \varepsilon_0 c \omega_L^2} I \longrightarrow \text{Low field seeker}$$



- Laguerre-Gauss beam
- Transverse confinement



- Crossed Gaussian beams
- Longitudinal confinement
- Tunable interatomic distance



$$\lambda=1\mu\text{m}, \Delta\theta=12^\circ$$

Intersite spacing = 5  $\mu\text{m}$

$$\omega_\perp \sim \omega_\parallel/2 = 12 \text{ kHz}$$

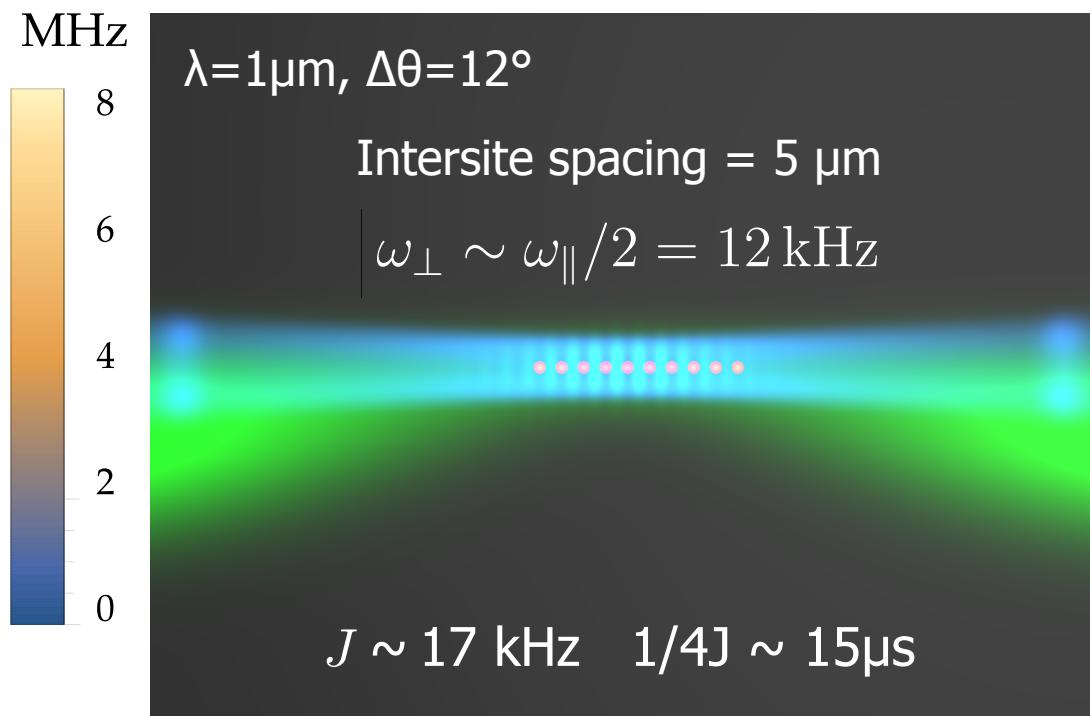
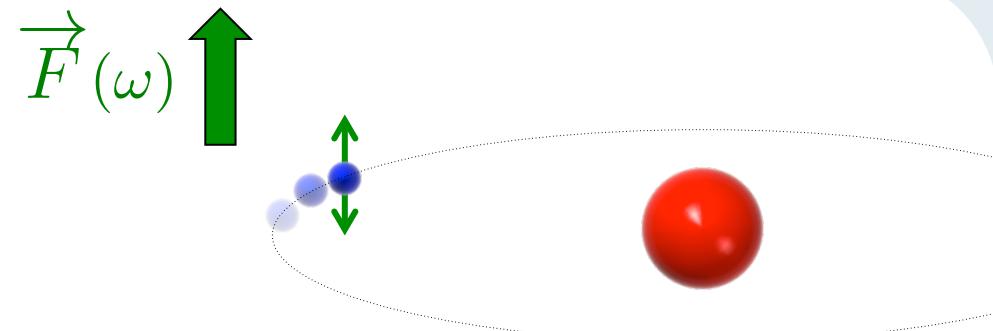
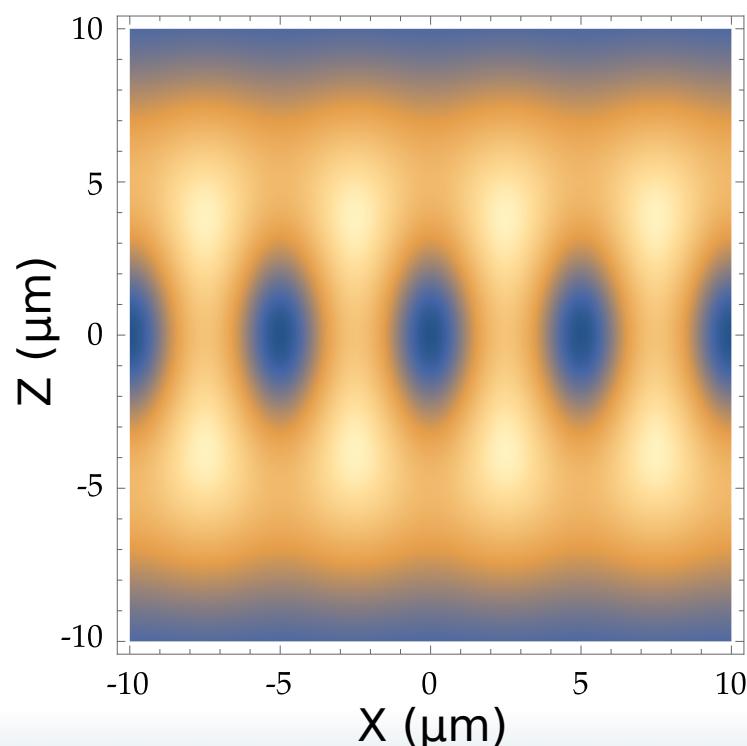
$$J \sim 17 \text{ kHz} \quad 1/4J \sim 15\mu\text{s}$$

# Trapping Rydberg atoms

## Ponderomotive potential

- Valence electron is almost free
- Positive ponderomotive energy

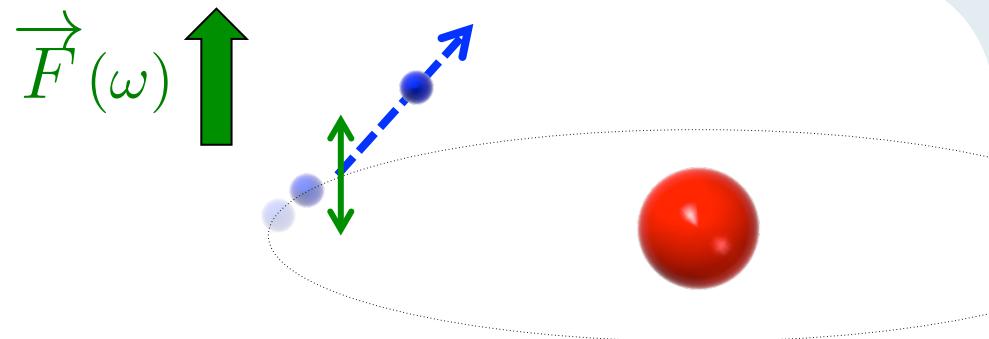
$$\mathcal{E} = \frac{e^2}{2m_e \varepsilon_0 c \omega_L^2} I \quad \rightarrow \text{Low field seeker}$$



# Trapping Rydberg atoms

## Ponderomotive potential

- Valence electron is almost free
- Positive ponderomotive energy
- MHz-deep lattices



## Photoionization?

- Detrimental for low- $l$  levels  
lifetimes  $\sim$  few 10  $\mu\text{s}$
- Negligible for circular levels!  
cross-section  $\sim 10^{-175} \text{ m}^2$

$$\lambda = 1\mu\text{m}, \Delta\theta = 12^\circ$$

Intersite spacing = 5  $\mu\text{m}$

$$\omega_\perp \sim \omega_\parallel/2 = 12 \text{ kHz}$$

$$J \sim 17 \text{ kHz} \quad 1/4J \sim 15\mu\text{s}$$

# Trapping Rydberg atoms

## Ponderomotive potential

- Valence electron is almost free
- Positive ponderomotive energy
- MHz-deep lattices
- No photoionization

## Coherence?

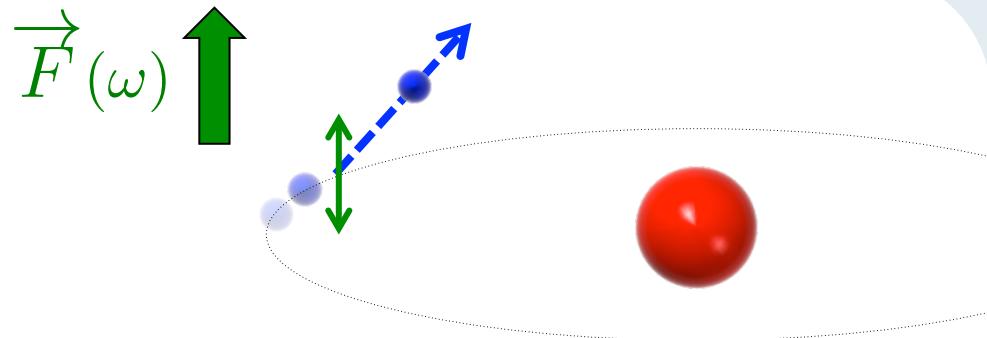
- Potential almost independent of the circular Rydberg level

Ground-state extension  $\sim 50$  nm

Electron-orbit radius  $\sim 250$  nm

→ Potential is averaged over the orbit

- Estimated coherence time  $\sim 0.2$  s



$$\lambda=1\mu\text{m}, \Delta\theta=12^\circ$$

Intersite spacing = 5  $\mu\text{m}$

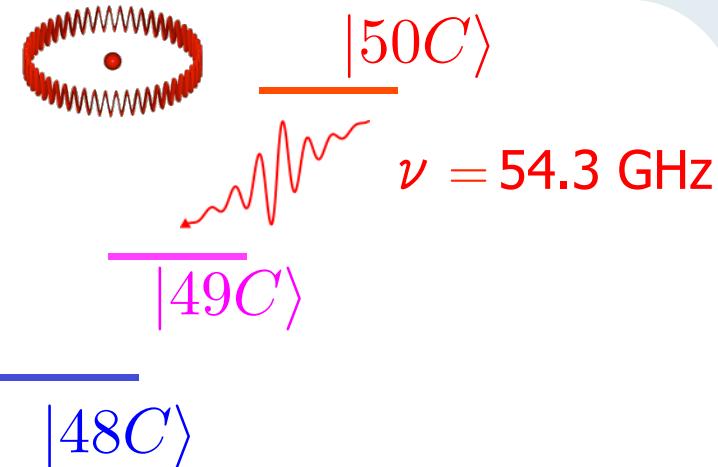
$$\omega_\perp \sim \omega_{||}/2 = 12 \text{ kHz}$$

$$J \sim 17 \text{ kHz} \quad 1/4J \sim 15\mu\text{s}$$

# Keeping Rydberg atoms

## Radiative decay

- Unique decay channel  $|nC\rangle \rightarrow |(n-1)C\rangle$
- Long natural lifetime  $\sim 30\text{ms}$   
... only 0.75ms for a 40-atom chain



$\lambda=1\mu\text{m}, \Delta\theta=12^\circ$

Intersite spacing = 5  $\mu\text{m}$

$$\omega_{\perp} \sim \omega_{\parallel}/2 = 12 \text{ kHz}$$

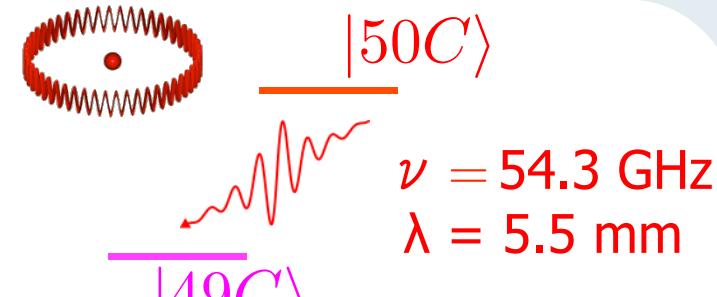


$$J \sim 17 \text{ kHz} \quad 1/4J \sim 15\mu\text{s}$$

# Keeping Rydberg atoms

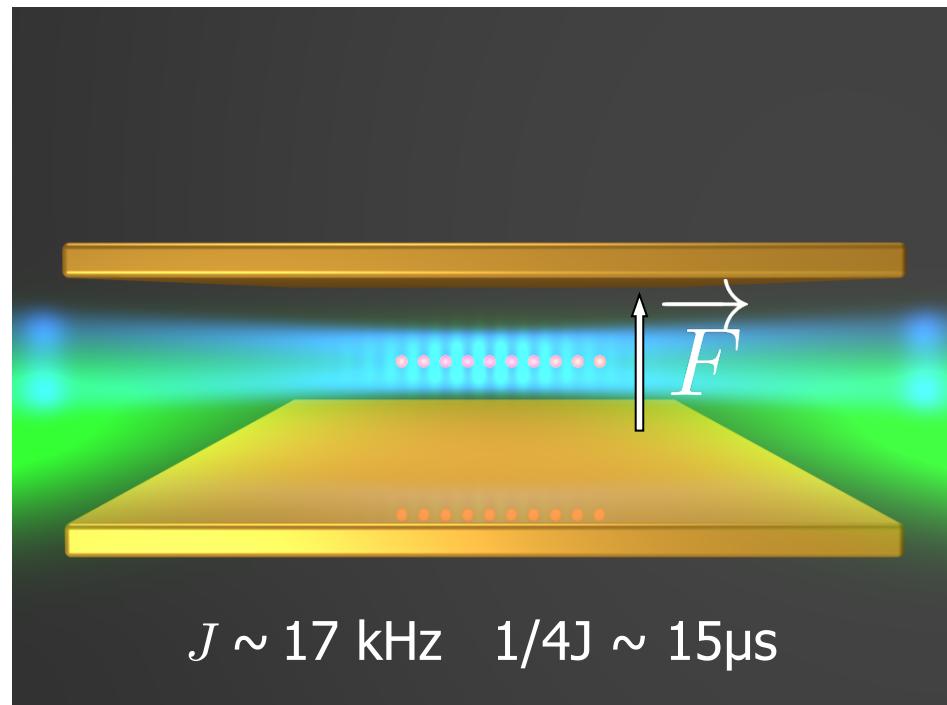
## Radiative decay

- Unique decay channel  $|nC\rangle \rightarrow |(n-1)C\rangle$
- Long natural lifetime  $\sim 30\text{ms}$   
... only 0.75ms for a 40-atom spin chain



## Inhibition of spontaneous emission

- Plane-parallel capacitor,  $D \leq \lambda/2$
- Emission of  $\sigma^+$  polarized photon is inhibited

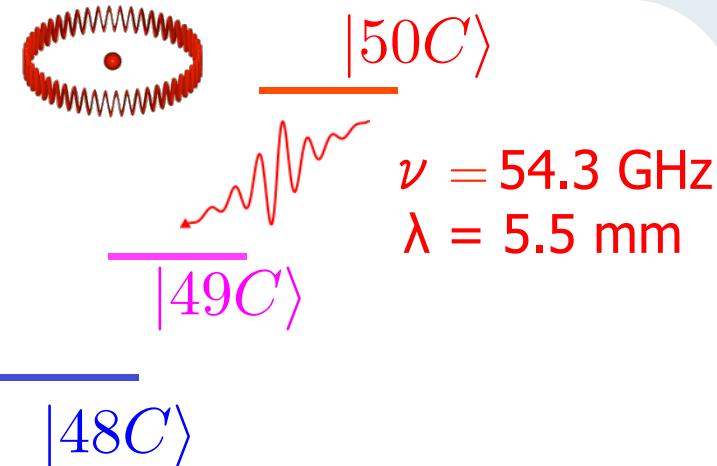


D. Kleppner [PRL 55, 2137 (1985)]

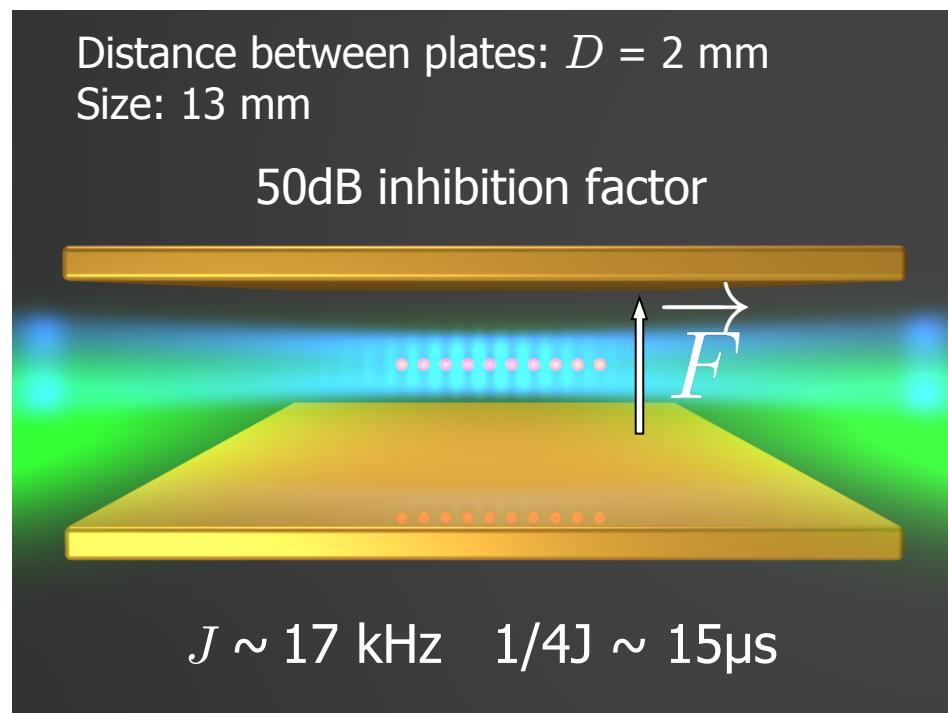
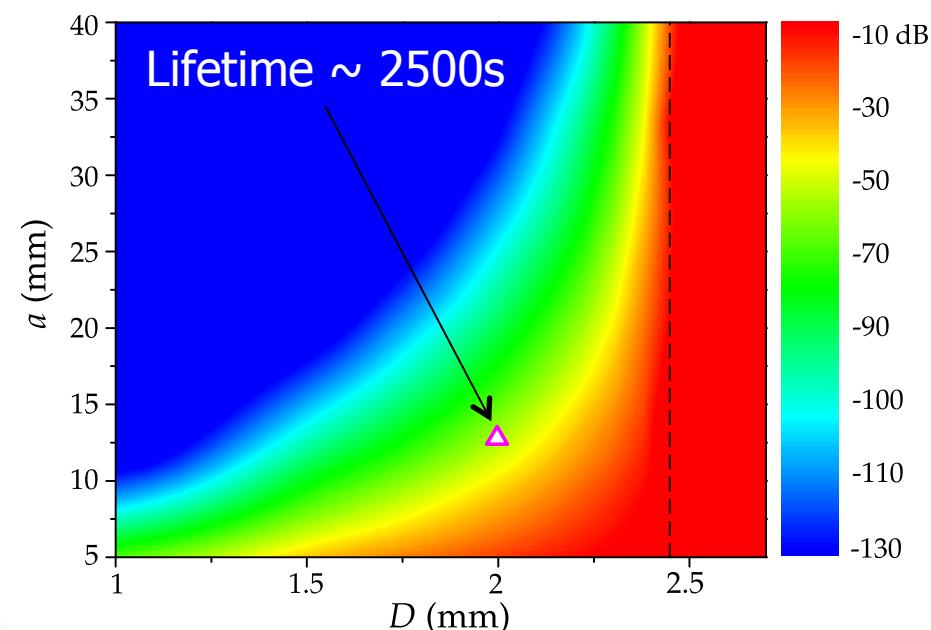
# Keeping Rydberg atoms

## Radiative decay

- Unique decay channel  $|nC\rangle \rightarrow |(n-1)C\rangle$
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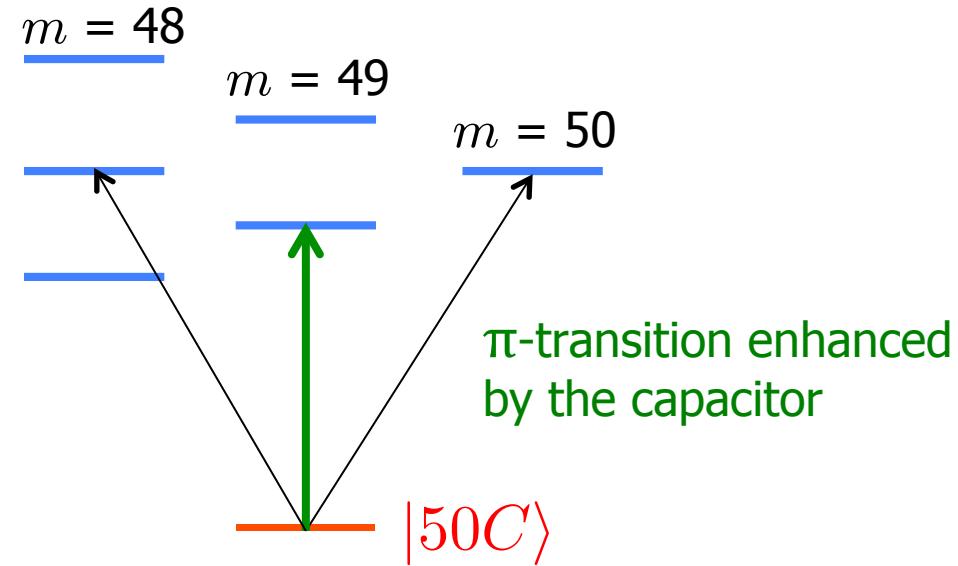
## Inhibition of spontaneous emission



# Keeping Rydberg atoms

## Limitations of the lifetime

- Background gas collisions:  $10^{-14}$  mbar required to reach 400s lifetime  
→ accessible in a cryostat environment
- Blackbody induced processes  
→ cryogenic temperature required:  $T \sim 0.5\text{K}$



# Keeping Rydberg atoms

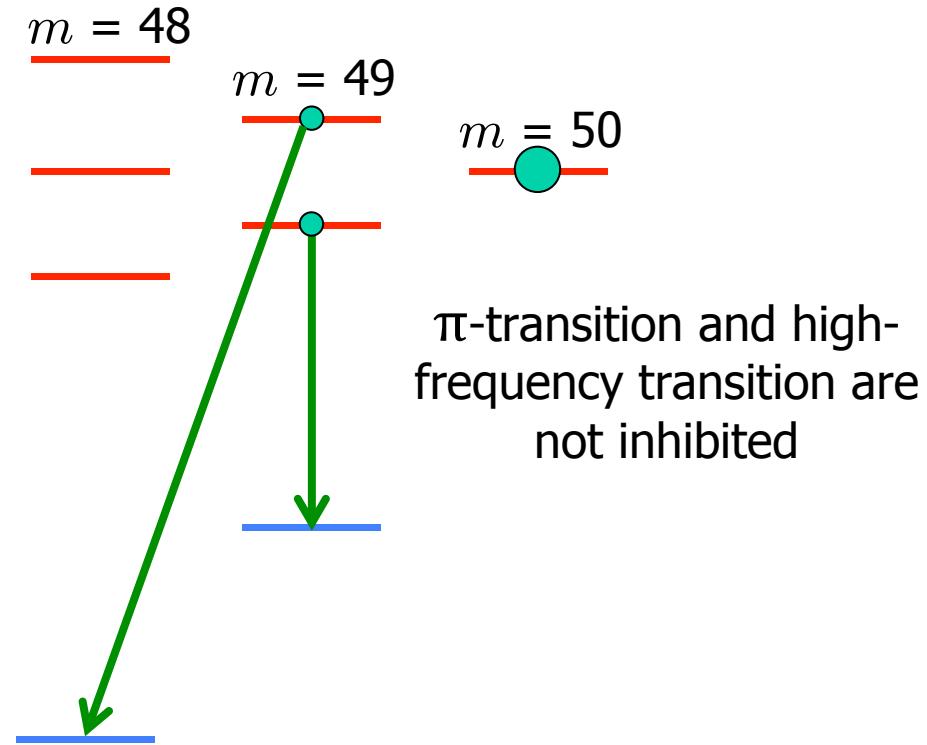
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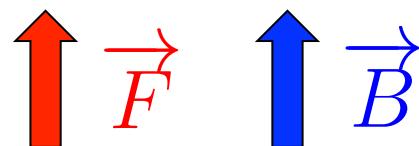
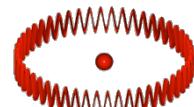
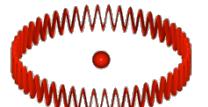
- Interaction-induced level mixing



# Keeping Rydberg atoms

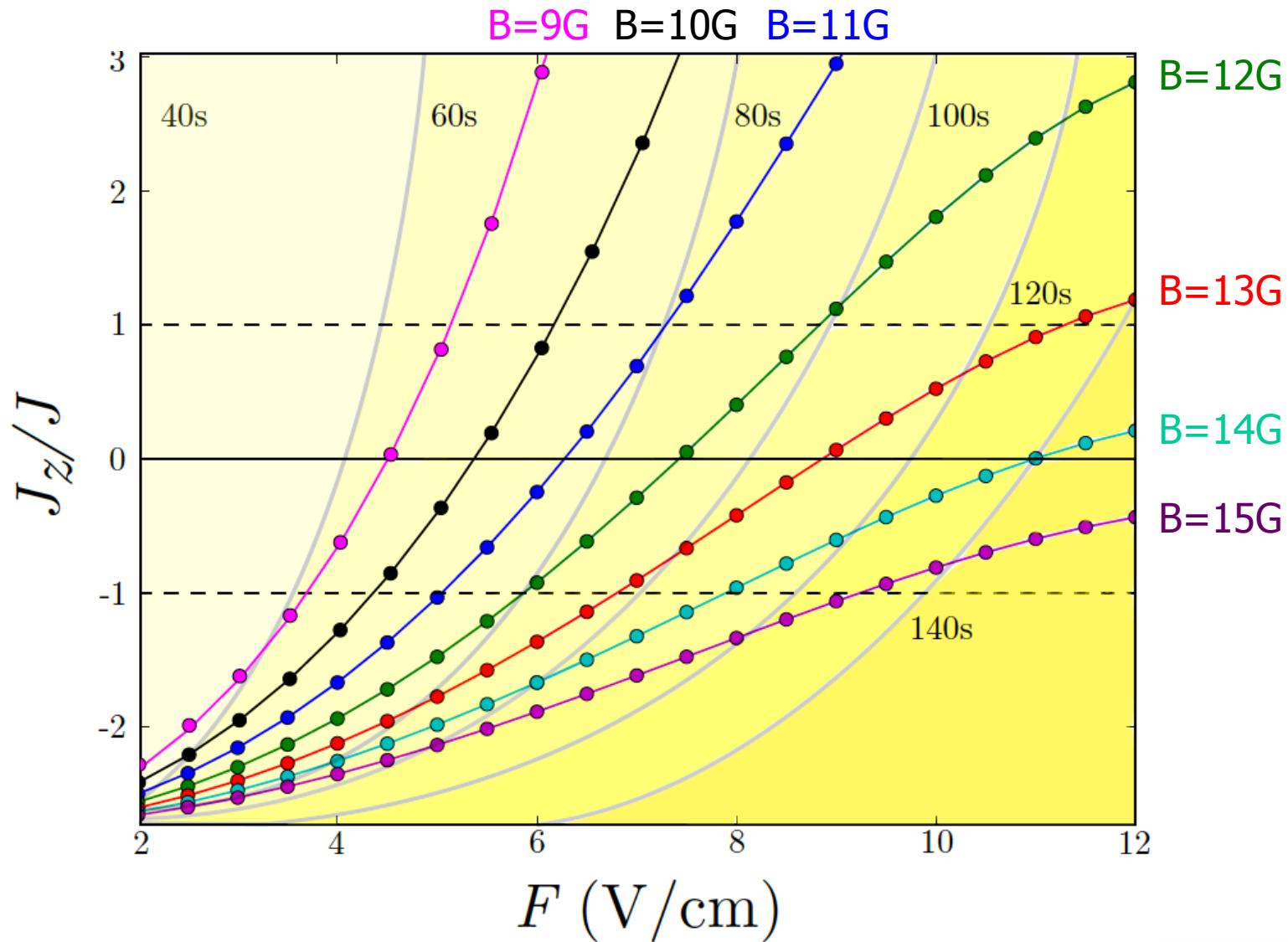
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- Background gas collisions:  $10^{-14}$  mbar required to reach 400s lifetime  
→ accessible in a cryostat environment
- Blackbody induced processes  
→ cryogenic temperature required:  $T \sim 0.5\text{K}$
- Interaction-induced level mixing → Big enough electric and magnetic fields are required



Competition with the tuning of the interaction?

# Keeping Rydberg atoms



# Keeping Rydberg atoms

## Limitations of the lifetime

- Background gas collisions:  $10^{-14}$  mbar required to reach 400s lifetime

→ accessible in a cryostat environment

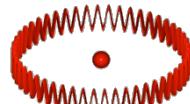
- Blackbody induced processes

→ cryogenic temperature  
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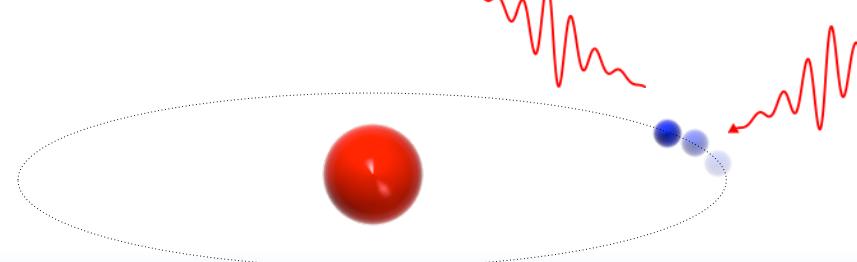
- Interaction-induced level mixing



Big enough electric and magnetic fields are required



- Compton elastic diffusion of trap photons?



# Keeping Rydberg atoms

## Summary

Cause	Lifetime (s)	
Residual spontaneous emission	2500	13mm square plates separated by 2mm
Blackbody induced processes	630	$T \sim 0.5\text{K}$
Level mixing	88	$B > 9\text{G} / F > 2\text{V/cm}$
Photoionization	$\infty$	
Collisions with background gas	400	$P \sim 10^{-14} \text{ mbar}$
Compton elastic diffusion	$> 180$	$P = 0.5\text{W}$
<b>Predicted lifetime</b>	<b>47</b>	

with  $J \sim 17 \text{ kHz}$   $1/4J \sim 15\mu\text{s}$

Single atom lifetime  $\sim 50 \text{ s}$

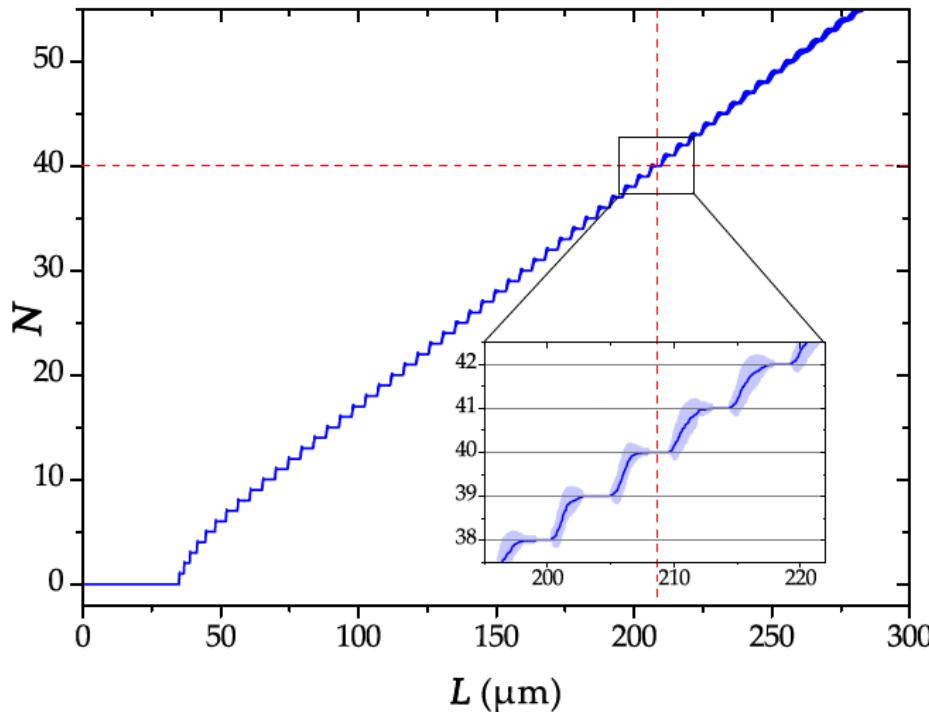
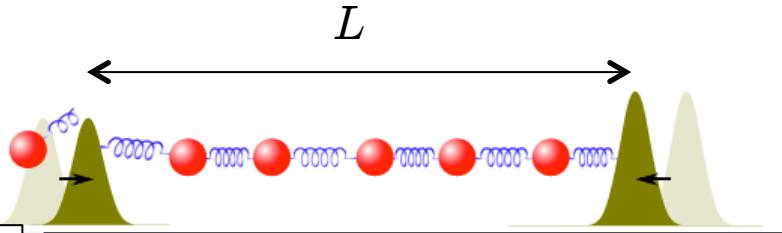
40-atom chain lifetime  $\geq 1 \text{ s}$

→  $\sim 10^5$  interaction cycles!

# Chain preparation

## Deterministic chain evaporation

- Cloud of cold atoms, near quantum degeneracy
- Excitation of Rydberg states and “circularization” inside the capacitor
- Evaporation of the chain

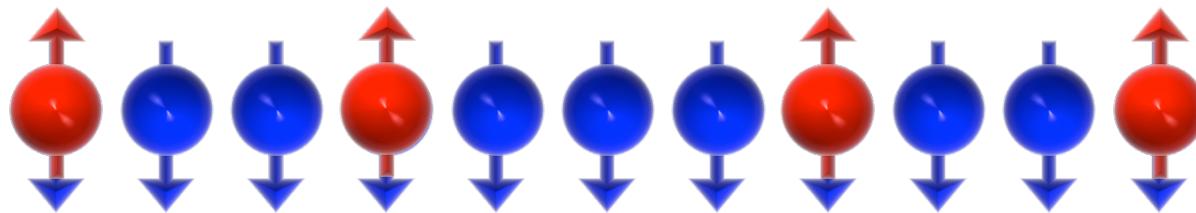


## Properties

- 40-atom chain prepared in  $\sim 1\text{s}$
- Variant of **evaporative cooling**:
  - 4 quanta of vibration
  - 100 nm position dispersion

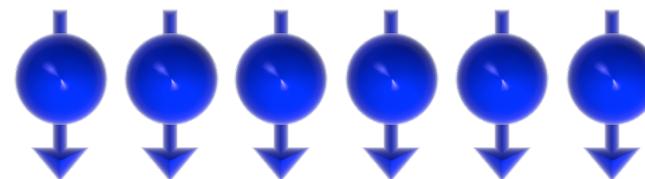
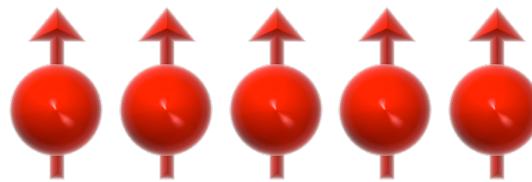
# Example of measurements

- Chain **initialization**: all atoms in the same spin state
- **Flipping** some atoms: atoms at the end of the chain (with single neighbour)

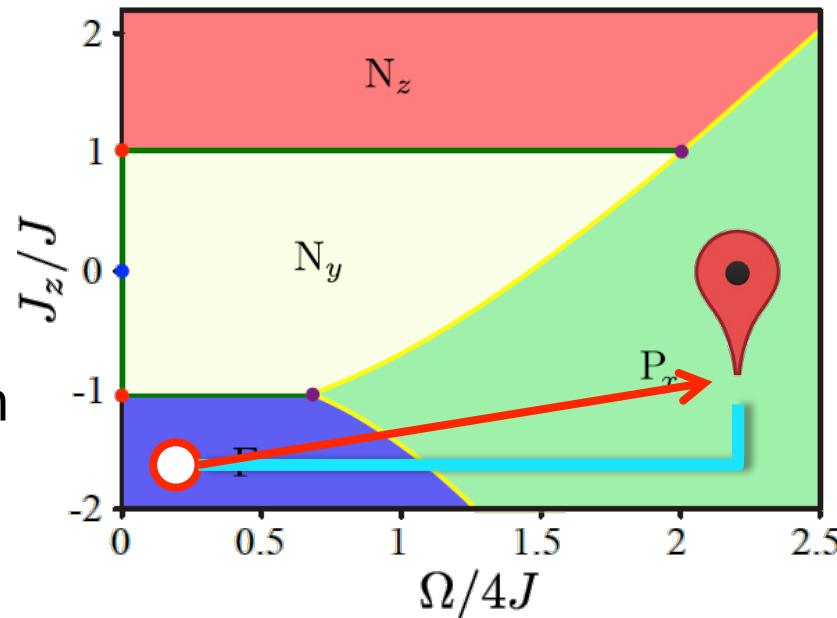


# Example of measurements

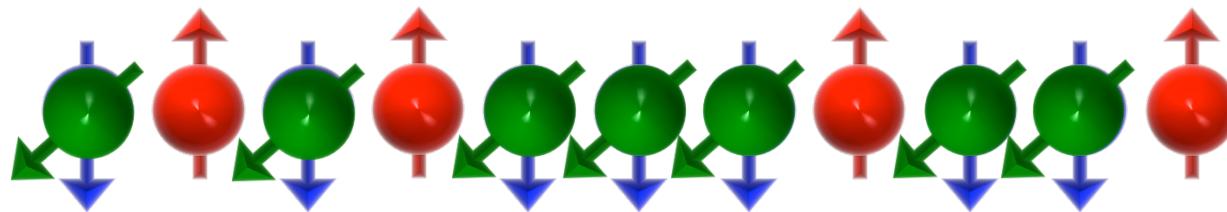
- Chain **initialization**: all atoms in the same spin state
- **Flipping** some atoms: atoms at the end of the chain (with single neighbour)



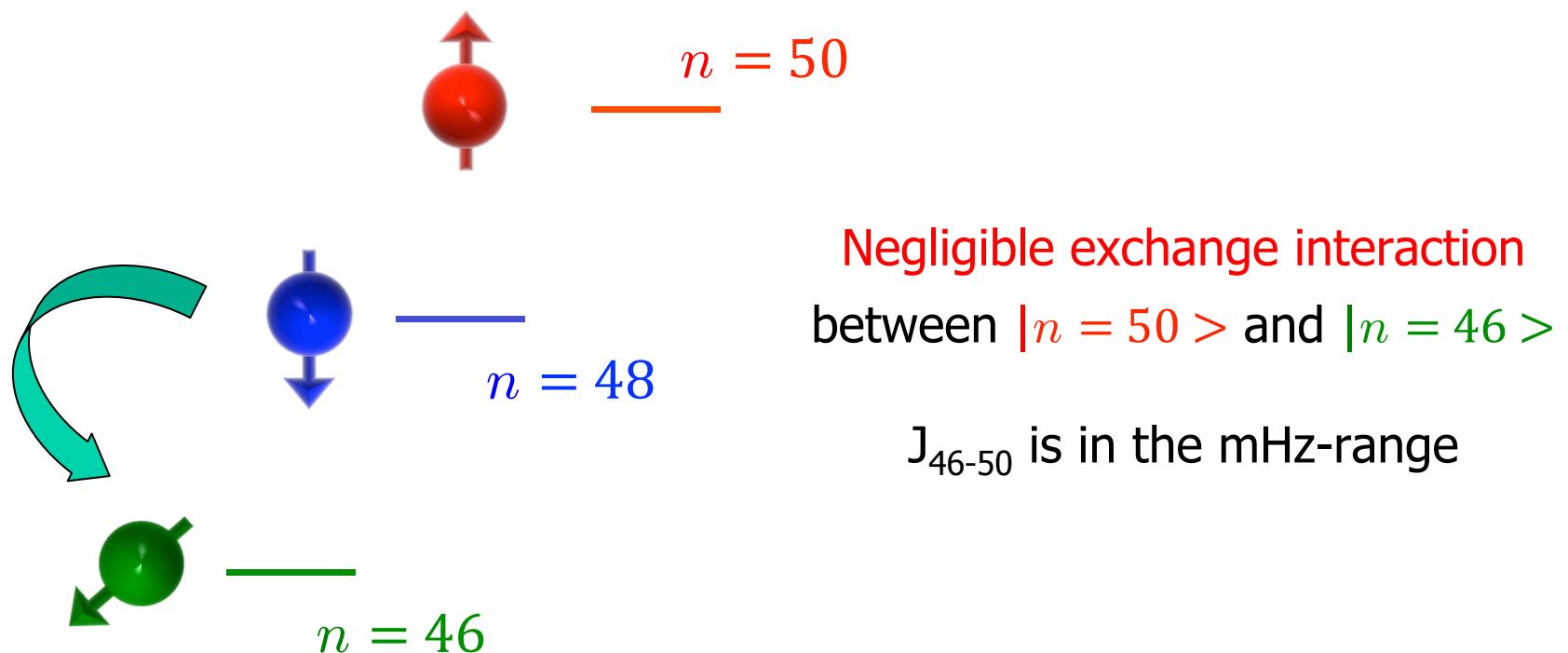
- Moving through the phase diagram:
  - Initial state in the ground state of the ferromagnetic phase
  - **Adiabatic evolution**: reconstruction of the phase diagram
  - **Quenches**: return to equilibrium, excitations...



# Example of measurements



- At the end of the sequence: **freeze the spin dynamics**

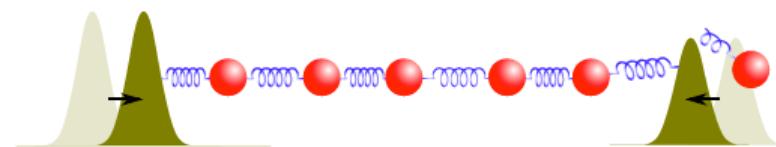


# Example of measurements

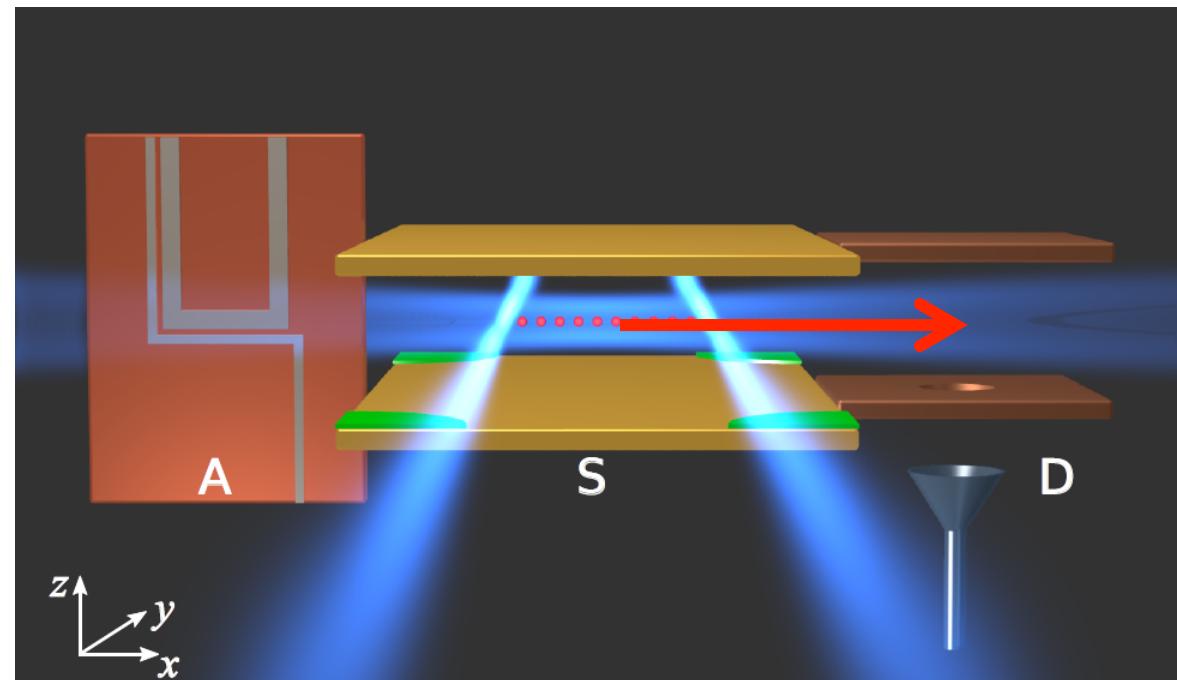
## Measurement of spin observables

- State-selective field ionization  $\longrightarrow$  Measurement of  $\sigma^z$
- Evaporation procedure is resumed

$\longrightarrow$  All atoms are detected  
one after the other



- Measurement of any spin component up to a global rotation of the spins
- Measurement of all spin correlations between every atoms



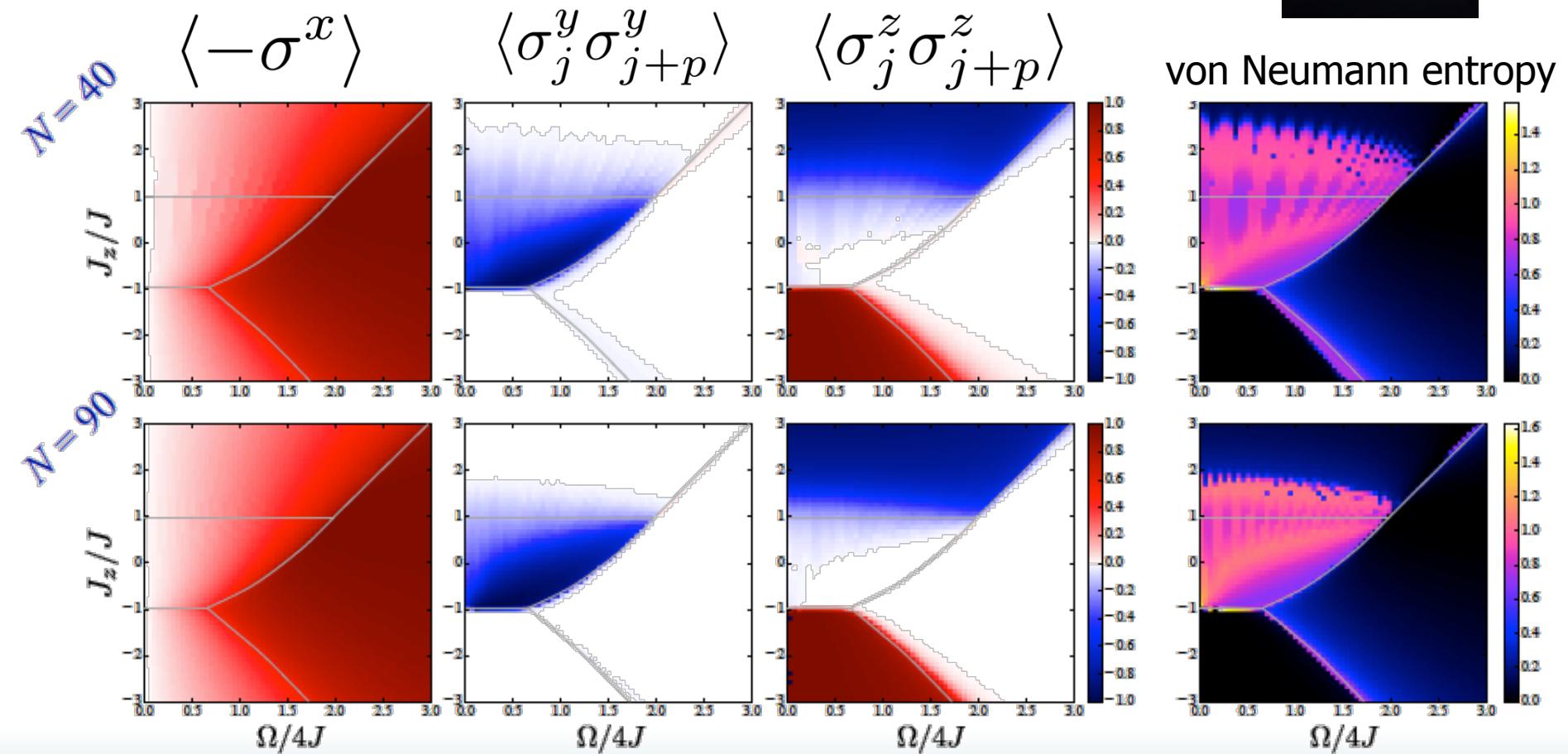
# Reconstruction of the phase diagram

## Numerical simulations

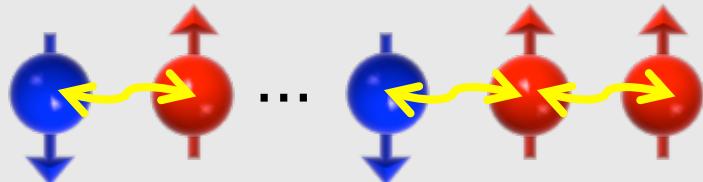
G. Roux, LPTMS (Orsay)



- Spin populations and correlations reveal the **phase transitions**
- Limited finite-size effects already with N=40 atoms



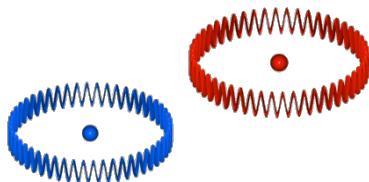
# Summary



Simulation of a chain of interacting spins-1/2

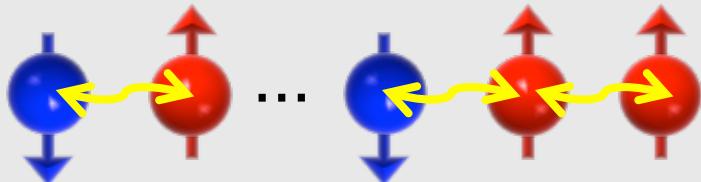
## Circular Rydberg atom quantum simulator

- Spin 1/2
- Defect free chain of 40 spins
- Long chain lifetime ( $\sim 1\text{s}$ ) and strong nearest-neighbour interactions ( $\sim 15\mu\text{s}$ )
  - Observation of  $10^5$  interaction cycles
- Fully tunable XXZ Hamiltonian



T. L. Nguyen et al., arXiv:1707.04397

# Summary



Simulation of a chain of interacting spins-1/2

## Circular Rydberg atom quantum simulator

- Exploration of long-term dynamics: return to equilibrium after a quench, quantum thermodynamics
- Effects of disorder
- Coupling to a common bosonic bath (motion)
- High frequency modulation: Floquet engineering
- Extension to 2D-protocols (e.g., spin-1 physics)

T. L. Nguyen et al., arXiv:1707.04397

# Thank you!



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